

# What can we learn from an improved search for free $N$ - $N$ -bar Oscillation ?

R. N. Mohapatra



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# Two ways to probe physics beyond standard model

- Direct search e.g in colliders such as LHC.
- Indirectly by searching for rare processes forbidden in SM.

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{M} \mathcal{O}^5 + \frac{1}{M^2} \mathcal{O}^6 + \dots + \frac{1}{M^n} \mathcal{O}^{4+n}$$

- One such rare effect is the neutrino mass,  $(LH)^2$  -d=5  
M not known yet. **Has B-L=2**
- Another class of processes- those violating baryon number: p-decay, NNbar osc.



# What new physics they reveal

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(i) d=6 operators (Weinberg; Wilczek, Zee'79)  $B-L=0$

$$\begin{aligned}\mathcal{O}_1 &= (d^c u^c)^* (Q_i L_j) \epsilon_{ij}, \quad \mathcal{O}_2 = (Q_i Q_j) (u^c e^c)^* \epsilon_{ij}, \quad \mathcal{O}_3 = (Q_i Q_j) (Q_k L_l) \epsilon_{ij} \epsilon_{kl} \\ \mathcal{O}_4 &= (Q_i Q_j) (Q_k L_l) (\vec{\tau} \epsilon)_{ij} \cdot (\vec{\tau} \epsilon)_{kl}, \quad \mathcal{O}_5 = (d^c u^c)^* (u^c e^c)^* .\end{aligned}$$

- Leads to  $p \rightarrow e^+ + \pi^0; p \rightarrow K^+ \bar{\nu}$  canonical GUT modes
- Current data bound the mass scale  $M > 10^{15}$  GeV
- They arise in GUT models with predictions (with uncertainties) and has spurred experimentalists to look for them for the last 30 years.

# B-L=2 B-violation

(ii) d=7 : Have B-L=2

(Weinberg; Weldon, Zee'80)

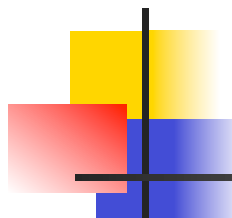
$$\begin{aligned}
 \tilde{\mathcal{O}}_1 &= (d^c u^c)^* (d^c L_i)^* H_j^* \epsilon_{ij}, & \tilde{\mathcal{O}}_2 &= (d^c d^c)^* (u^c L_i)^* H_j^* \epsilon_{ij}, \\
 \tilde{\mathcal{O}}_3 &= (Q_i Q_j) (d^c L_k)^* H_l^* \epsilon_{ij} \epsilon_{kl}, & \tilde{\mathcal{O}}_4 &= (Q_i Q_j) (d^c L_k)^* H_l^* (\vec{\tau} \epsilon)_{ij} \cdot (\vec{\tau} \epsilon)_{kl}, \\
 \tilde{\mathcal{O}}_5 &= (Q_i e^c) (d^c d^c)^* H_i^*, & \tilde{\mathcal{O}}_6 &= (d^c d^c)^* (d^c L_i)^* H_i, \\
 \tilde{\mathcal{O}}_7 &= (d^c D_\mu d^c)^* (\bar{L}_i \gamma^\mu Q_i), & \tilde{\mathcal{O}}_8 &= (d^c D_\mu L_i)^* (\bar{d}^c \gamma^\mu Q_i), \\
 \tilde{\mathcal{O}}_9 &= (d^c D_\mu d^c)^* (\bar{d}^c \gamma^\mu e^c).
 \end{aligned}$$

$$\rightarrow n \rightarrow e^- \pi^+$$

(iii) d=9:  $u^c d^c d^c u^c d^c d^c$  B-L=2  $\rightarrow n - \bar{n}$

(RNM, Marshak'80; Glashow'80)

- Curious feature: odd d B-L=2 and even d, B-L=0; No B-L=1 due to Lorentz invariance



# Neutrino mass- $NN\bar{}$ connection

- Both have odd  $d$  and break  $B-L$  by 2 units;
  - If neutrino is Majorana, it breaks  $\mathcal{L}$ -part of  $\mathbf{B-L}$  by  $2_{\text{units}}$
  - $N-N\bar{}$  oscillation breaks  $B$ -part of  $B-L$  and provides complementary information on  $\nu$  mass physics
  - Are there models where this happens ?
  - Yes.  $NN\bar{}$  is a generic prediction of theories where seesaw mechanism for  $\nu$  mass is embedded into quark-lepton unification as in grand unified theories
- Seesaw breaks  $\mathcal{L}$  and  $QL$  unification converts to  $B=2$  !*



# Today's plan

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- Upper bound on  $N_{\text{Nbar}}$  transition time  
Q-L unified 224 model
- Comments on GUT models and  $N_{\text{Nbar}}$

# SEESAW+QUARK-LEPTON UNIFICATION and NN-bar

- Seesaw model for neutrino masses:
- Standard model + RH neutrino  $N$
- Two sets of Higgs fields:
  - (i) SM doublet for usual fermion masses:  $\phi$
  - (ii) New Higgs field to give mass to RH neutrino:  $\Delta_R^0$
- $\langle \Delta^0 \rangle_R \neq 0, \langle \phi \rangle \neq 0$  .

$$(\nu, N) \begin{pmatrix} 0 & h \langle \phi \rangle \\ h \langle \phi \rangle & f \langle \Delta_R \rangle \end{pmatrix} \begin{pmatrix} \nu \\ N \end{pmatrix}$$
- $m_\nu \approx (m_e)^2 / M_R \ll m_e$  (Seesaw formula)

# A UNIFIED TEV SCALE

## EMBEDDING OF SEESAW

If Q-L unified at the seesaw, a model is

$$SU(2)_L \times SU(2)_R \times SU(4)_C \begin{pmatrix} u & u & u & \nu \\ d & d & d & e \end{pmatrix}_{L,R}$$

→ SU(4) generalization of the seesaw Higgs field  $\Delta_R$  has partners  $\Delta_{qq}$  connecting to quarks

→ N-N-bar Feynman graph;

(Mohapatra, Marshak'80)

→ No proton decay.

$$\tau_{n-\bar{n}} = \hbar / \delta m_{n-\bar{n}} \sim M^5 / \Lambda^6$$



$$\frac{\lambda f^3 \nu_{BL}}{M_{\Delta}^6}$$

Colored seesaw partners at TeV scale →  $\tau \sim 10^{10-11}$  sec.





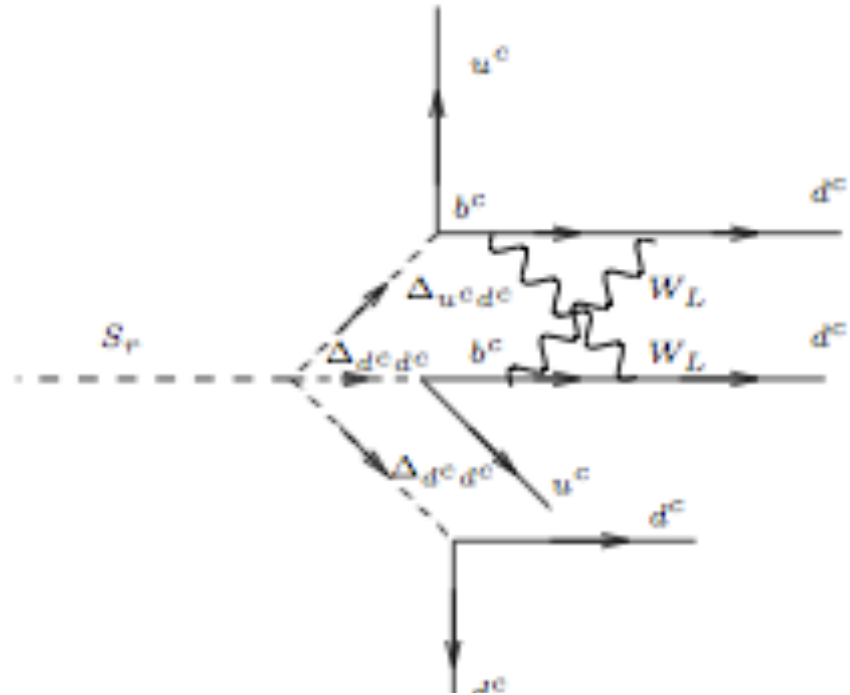
# Two consequences of $SU(4)_c$

- Couplings of  $\Delta_{qq}$  elements of neutrino mass matrix !!
- They lead to FCNC effects and thus constrained !!
- Is there a solution ? (Babu, Dev, RNM'09; Fortes, Babu, RNM (to appear))
- Yes, inverted hierarchy for nu masses with type II seesaw:  
■ Fixes coupling matrix:  $f_{dd} = \begin{pmatrix} 0 & 0.95 & 1 \\ 0.95 & 0 & 0.01 \\ 1 & 0.01 & -0.063 \end{pmatrix} c$
- Masses  $M_{\Delta_{dd}} \sim 10 TeV$ ;  $M_{\Delta_{ud}} \sim 1 TeV$   
satisfy all FCNC constraints
- NNbar osc loop effect; without  $SU(4)$ ,  $f_{11}$  arbitrary as is NNbar

# Neutron-anti-neutron oscillation

- Nnbar osc. Is a loop effect:

$$F_{ud} = U_{CKM} f_{dd}$$



$$G_{n-\bar{n}} \simeq \frac{f_{ud,11} f_{ud,13} f_{dd,13} \lambda v_{BL}}{M_{\Delta_{u^c d^c}}^4 M_{\Delta_{d^c d^c}}^2} \frac{g^4 V_{td}^2 m_b^2 m_t^2}{(16\pi^2)^2 m_W^4} \log\left(\frac{m_b^2}{m_W^2}\right)$$



# Requirement of Baryogenesis

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- Basic message:

(i) If all particles involved in  $NN\bar{b}$  diagram are  $\sim \text{TeV}$  mass, baryogenesis must be below electroweak phase transition temp. (PSB) (Babu, Nasri, RNM'07)

(ii) Cosmology of baryogen. predicts a lower bound for

$$G_{n-\bar{n}}$$

$$G_{nn} > 10^{-31} \text{ GeV}^{-5} \rightarrow \tau_{n\bar{n}} \leq 10^{11} \text{ sec}$$

(Without cosmology restrictions, scale  $c$  of the  $f$ -couplings is unknown and  $NN\bar{b}$  coupling could be anything; cosmology puts a lower bound on  $C \sim .7$ ):



# Basic outline of baryogenesis

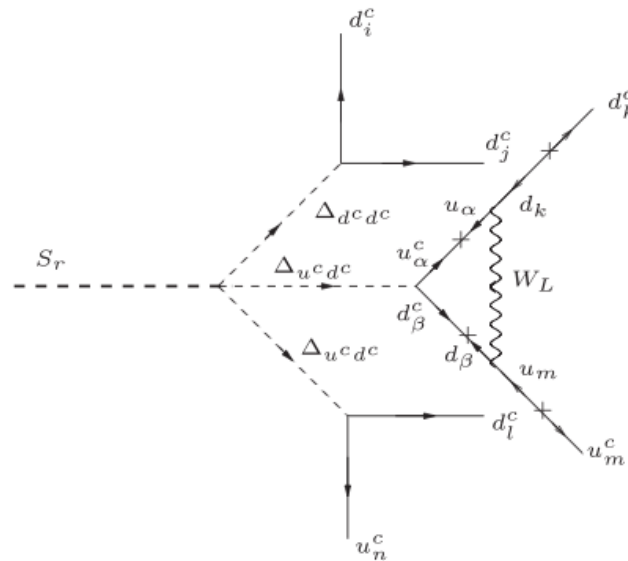
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- Above EWPT  $\rightarrow T \sim v_{wk}$ , both B+L and B=2 violating processes in eq.  $\rightarrow$  all asymmetry erased. Need a fresh start
- Means baryogenesis must occur below EWPT
- $\rightarrow$  particle decay  $T_d$  below EWPT and above .1 GeV

$$T_d \simeq \left[ \frac{36\lambda^2 (\text{Tr}[f^\dagger f])^3 M_{\text{Pl}} M_S^{13}}{(2\pi)^9 1.66 g_*^{1/2} (6M_\Delta)^{12}} \right]^{1/2}$$
$$\simeq 6.1 \text{ GeV}^{1/2} \left( \frac{M_S^{13}}{M_\Delta^{12}} \right)^{1/2}.$$

# Baryogenesis diagram

- Re  $\Delta_R$  decay to 6 quarks source of baryons.

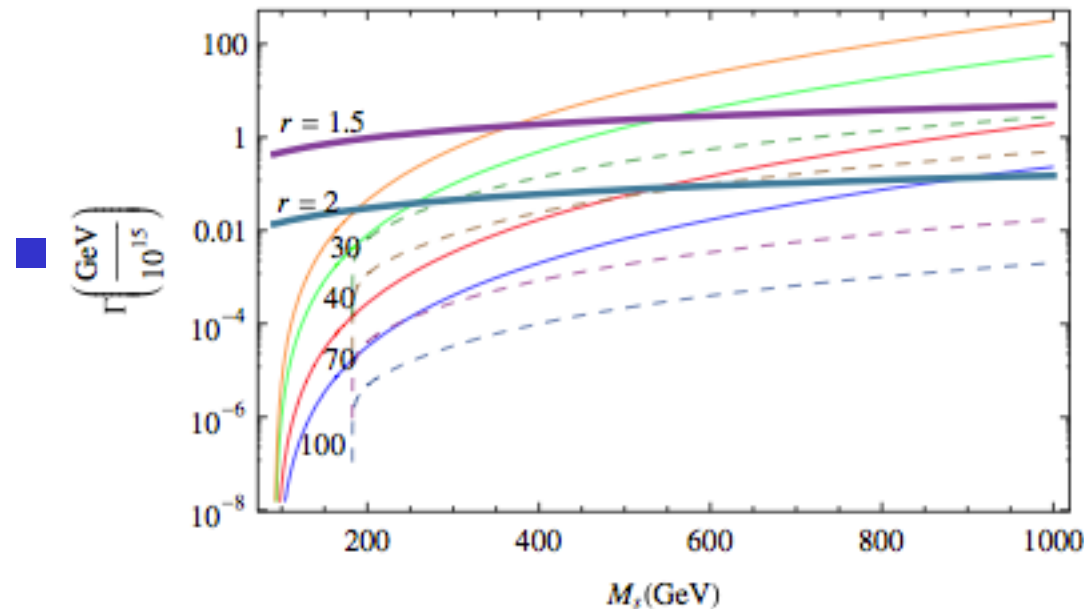


- $\rightarrow$  6q decay must dominate over others.

# Six quark decay dominance constraint

- $$\Gamma_{Z f \bar{f}} \sim \frac{.07 M_S^5}{M_{Z'}^6} \text{GeV}$$

$$\Gamma(S_r \rightarrow 6q^c) \simeq \frac{36}{(2\pi)^9} \frac{(\text{Tr}[f^\dagger f])^3 \lambda^2 M_S^{13}}{(6M_\Delta)^{12}},$$



$$v_{\text{BL}} < 100 \text{ TeV}$$



# Summary of constraints

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- Low scale baryogenesis in  $SU(2)_L \times SU(2)_R \times SU(4)_C \rightarrow$ 
  - i)  $M_{\Delta_{qq}} > M_S$
  - ii)  $1 \text{ GeV} < T_{S-\text{decay}} < 100 \text{ GeV}$
  - iii)  $\Gamma_{S \rightarrow 6q} > \Gamma_{S \rightarrow Z q \bar{q}}$
  - iv) A neutrino mass fit+FCNC constraints
- These constraints upper bound NNbar transition time
$$\begin{aligned} v_{BL} &< 100 \text{ TeV} & \tau &< 10^{11} \text{ sec.} \\ &< 30 \text{ TeV} & \tau &< 10^{10} \text{ sec.} \end{aligned}$$



# Color sextet scalars and LHC connection

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- FCNC, Baryogenesis  $\rightarrow$  color sextet spectrum:

$$M_{\Delta_{ud}} < M_{\Delta_{dd}} \ll M_{\Delta_{uu}}$$

- With  $M_{\Delta_{ud}} \leq TeV$
- Couples to  $ud, dt, bu, \dots$
- Should show up at LHC



# Production and decay

- TeVColor sextets are an inherent part of both models ;  
Can be searched at LHC:

(I) **Single production:**  $ud \rightarrow \Delta_{ud} \rightarrow tj$

xsection calculated in (RNM, Okada, Yu' 07;) resonance peaks above SM background- decay to tj; B(tj) suppressed in this model.

(II) **Drell-Yan pair production**  $q\bar{q} \rightarrow G \rightarrow \Delta_{ud} \bar{\Delta}_{ud}$

- Leads to  $tj\bar{t}\bar{j}$  final states: **LHC reach  $\sim$  TeV**

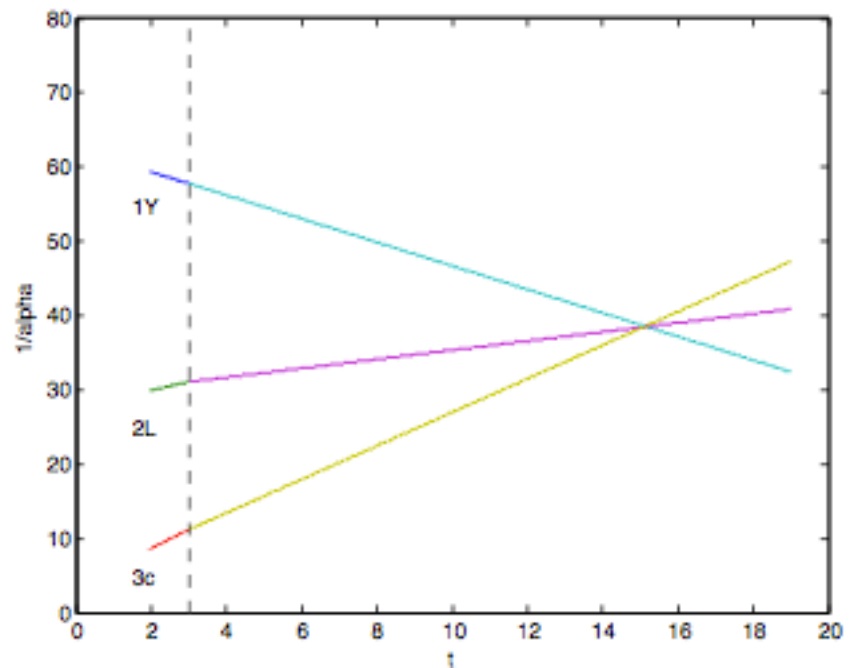
(Chen, Rentala, Wang; Berger, Cao, Chen, Shaughnessy, Zhang' 10; Han, Lewis' 09)

# Non-SUSY SO(10) – Another predictive model for NNbar

- Coupling unification fixes the mass scales as in the case of proton decay:
- In a minimal SO(10) embedding of seesaw,  $f_{ab}$  determined from fermion mass fits
  - (Babu, Mohapatra'93; Fukuyama, Okada'02; Bajc, Senjanovic, Vissani'02; Goh, Mohapatra, Ng'03 Babu, Macesanu'05; Bertolini, Malinsky, Schwetz'06; Joshipura, Patel'11)
- Predicts correct  $\theta_{23}, \theta_{12}$  and  $\sin^2 2\theta_{13} \simeq 0.09$
- Model has diquarks at sub-TeV scale to have unification and they lead to observable NNbar !

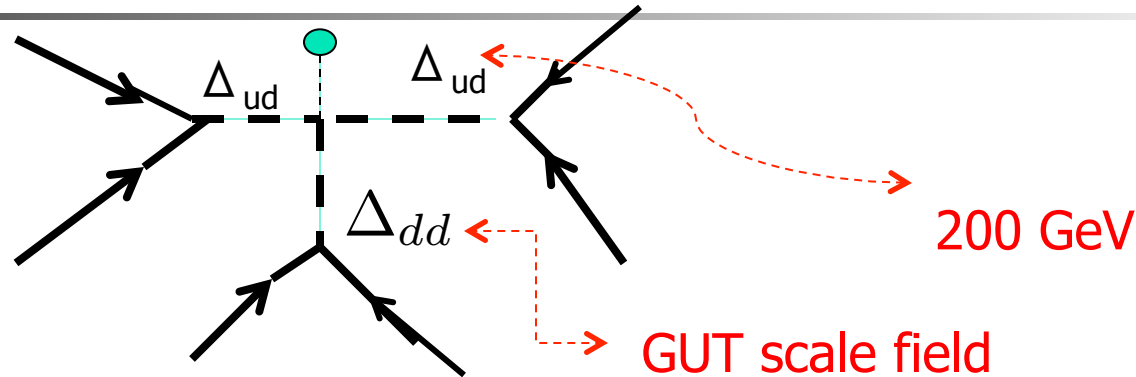
# New Unification profile and NNbar in predictive SO(10)

- Non-SUSY SO(10) does not unify without low scale particles,
- Coupling unif with sub-TeV  $\Delta_{ud}(6, 1, \frac{1}{3})$  + 2 SM triplets;
- Predicts seesaw scale near  $M_U \sim 10^{16}$  GeV;
- $\Delta_{ud}$  mass  $\sim 2$  TeV
- $M_U \sim 10^{15.7}$  GeV  $\rightarrow$   
 $\tau_{p \rightarrow e^+ + \pi^0} \simeq 3.2 \times 10^{34} \text{ yrs}$  close to current limit.



# Estimate of N-N-bar oscillation time

Diagram:



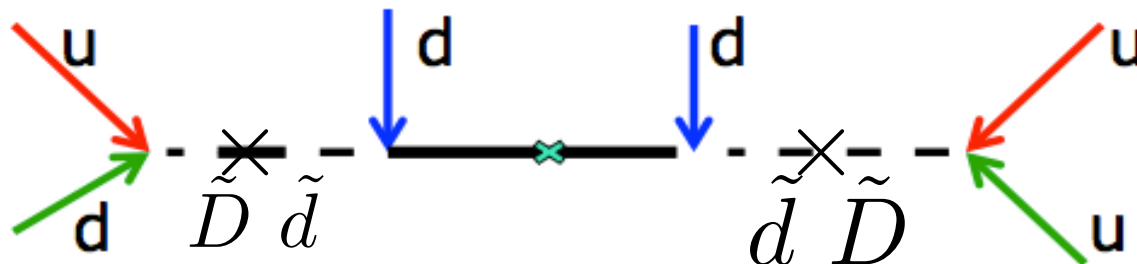
$$G_{\Delta B=2} \simeq \frac{\lambda f_{11}^3 \eta^3}{\lambda' M_U M_{\Delta_{ud}}^4} \simeq \frac{\lambda}{\lambda'} 10^{-33} \text{GeV}^{-5}$$

Predicts  $\tau_{n-\bar{n}} \sim 10^{10} - 10^{13} \text{ sec.}$

Constraints of baryogenesis reduces this by two orders of magnitude. (Babu's talk)

# A SUSY $E_6$ Model for $NN\bar{b}$

- $E_6$  contains  $SO(10)$  + extra vectorlike fermions:
- $\{27\} = \{16\} + \{10\} + \{1\}$
- Under  $SU(5)$ ,  $\{10\} = \{5\} + \{5^*\}$  ( $D, E^-, E^0$ )
- d-quark mixes with vectorlike D-quark and leads to  $NN\bar{b}$  oscillation operator: not connected to  $\nu$  mass: (RNM, Valle'86)

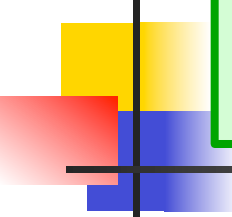




# Non-SUSY $E_6$ – No $NN\bar{b}$

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- Key to  $NN\bar{b}$  in  $SO(10)$ -  $\{126\}^4$  coupling needed for baryogenesis.
- $E_6$  contains  $SO(10) \times U(1)$  and this coupling is forbidden by the  $U(1)$ ; hence no  $NN\bar{b}$ .



# Benchmark goal for ruling out new physics scenarios

- **No  $N\bar{N}$  oscillation till  $\sim 10^{10} - 10^{11}$  sec.**

★ Will rule out a class of  $SU(2) \times SU(2) \times SU(4)_C$  models for post sphaleron baryogenesis for  $v_{BL} < 30-100$  TeV.

★ Will rule out a class of  $SO(10)$  models for neutrino masses that predicted recently observed large  $\theta_{13}$  if it is to explain the origin of matter.

**THANK YOU !!**

# Implications of NN-bar observation for other physics

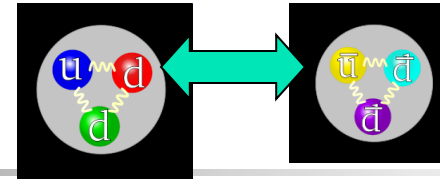
- Light color sextets  $\rightarrow$  FCNC effects **e.g.**  
D0 dimuon anomaly via B-B-bar mixing and B-decay CP asym. etc. (Babu, Fortes, RNM'11)
- EDM of neutron- two loop diagram
- Strange dibaryon decay:  $NN \rightarrow KK + X$  (Glashow)  
Mediated only by  $\Delta_{ud.dd}$  **Huge in our model**; recommend strange non-leptonic dibaryon decay mode search in p-decay!!



# What else can we learn from direct $N\bar{N}$ search?

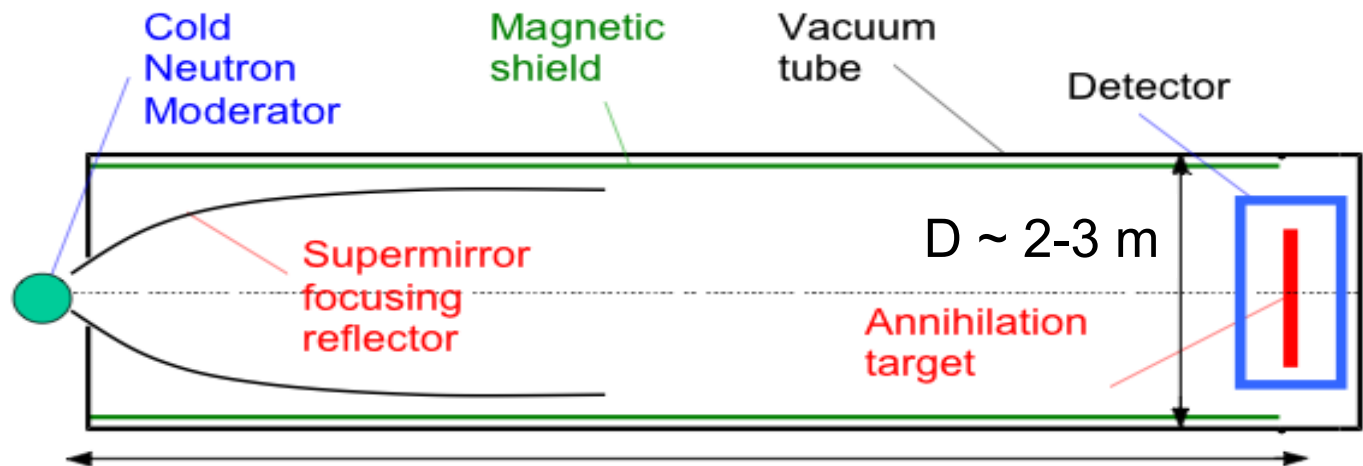
- Can test some dark matter hypothesis e.g. if a dark neutron  $n'$  is dark matter (ADM models):
- $n' \longleftrightarrow n$  oscillation can deplete dark matter density and this can be searched for in direct  $n\bar{n}$  searches; current limit  $> 1$  s. (Bento, Berezhiani)  
(possibly a signal ?)
- If  $N\bar{N}$  is discovered, it will put the strongest limit on CPT violation- (Okun; Addazzi, BLV2011)

# Search for N-N-bar Osc. current status



- **Free neutron oscillation in reactors: generic setup**

(talks by Snow, Young)



$L = 300 \text{ m}$

with  $L \sim 90 \text{ m}$  and  $\langle t \rangle = 0.11 \text{ sec}$   
measured  $P_{n\bar{n}} < 1.6 \times 10^{-18}$

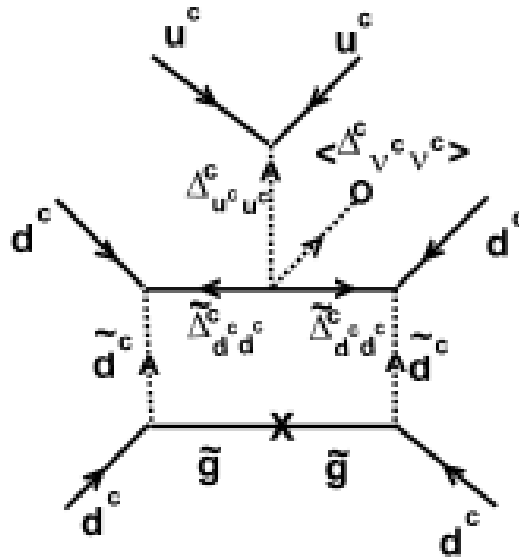
$\tau > 8.6 \times 10^7 \text{ sec}$

- **Current bound (ILL' 94)**

- **No new search after that**

# Estimate of N-N-bar with susy

## New Feynman diagram for N-N-bar osc.



$$G_{N-\bar{N}} \simeq \frac{f \tilde{m}}{\lambda^2 M_{\text{seesaw}}^2 v_{\text{uk}}^2}$$

$M_{\text{seesaw}} \sim 10^{11} \text{ GeV}$ , typical  $f, \lambda, \tau_{N-\bar{N}} \sim 10^{10} \text{ sec.}$

Observable N-N-bar osc for  **$M_{\text{seesaw}} \sim 10^{11} \text{ GeV}$** .

(Dutta, Mimura, RNM; PRL (2006))



# Scale reach of NNbar

**SM particles**  $O_{\Delta B=2} = \frac{1}{M^5} u^c d^c d^c u^c d^c d^c$  **d=9**

$$\delta m_{n-\bar{n}} = O_{\Delta B=2} \Lambda_{QCD}^6 \quad (\text{Lattice talks})$$

**Note  $M^5$  suppression**

$$\tau_{n-\bar{n}} = \hbar / \delta m_{n-\bar{n}} \sim M^5 / \Lambda^6 \rightarrow \tau_{n\bar{n}} \sim 10^8 s. \quad M \approx 10^{5.5} GeV$$

**TeV diquarks:**  $\rightarrow \Delta_{u^c d^c} \rightarrow \frac{1}{\Lambda \Lambda} d^c d^c \Delta_{u^c d^c} \Delta_{u^c d^c}$

$$M \simeq 10^{15} GeV$$

Same scale reach as proton decay but probes different physics, more closely related to nu mass

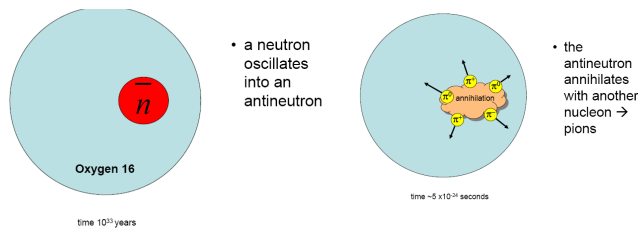


# Spin flip issue

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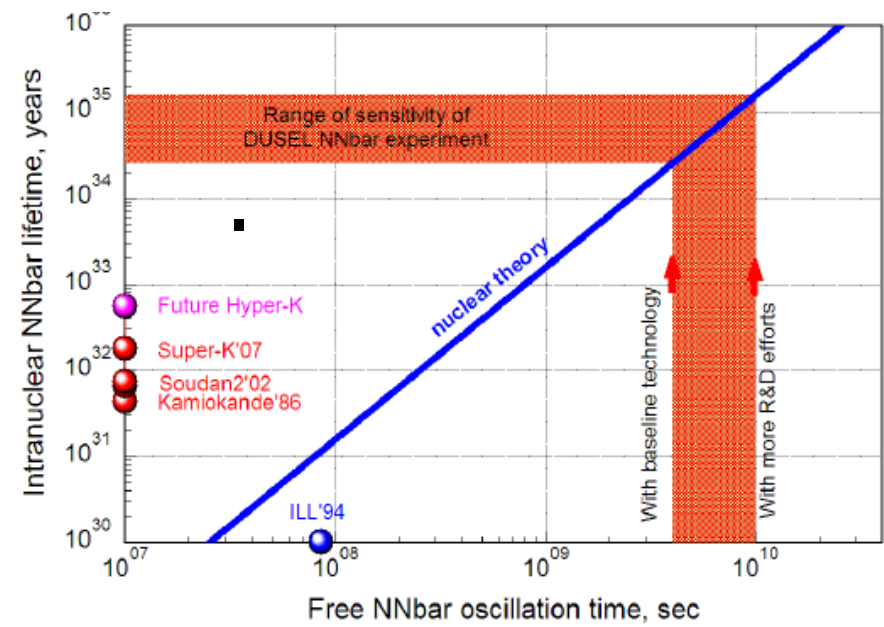
# Two ways to search for $\Delta B=2$ processes

- Free vs bound neutron oscillation: later  $\rightarrow (NN \rightarrow \pi' s)$



$$\tau_{Nuc} = R \tau_{free}^2$$

$$R = 0.3 \times 10^{23} \text{ sec}^{-1}$$



(Plot by Y. Kamyshev)

$$\tau_{n\bar{n}} > 2.44 \times 10^8 \text{ sec. (S-K, Abe et al.)}$$

- Free oscillation search much more effective !!